

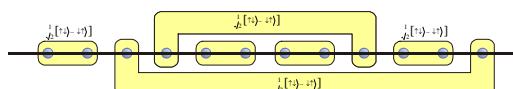
# Scaling and dynamics in quantum spin chains with bond randomness

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## Random Singlet model and universality

$$\hat{H} = \sum_j J_j \vec{S}_j \vec{S}_{j+1} \quad |J_j| > 0, \langle J_j \rangle = J$$

- The  $S=1/2$  spin chain is **critical** in the absence of disorder.
- Disorder is **relevant** in low-dimensional systems.
- Regardless of the details or strength of disorder, the low-energy properties of random-bond spin chains are as those in the **infinitely** disordered Random Singlet state:

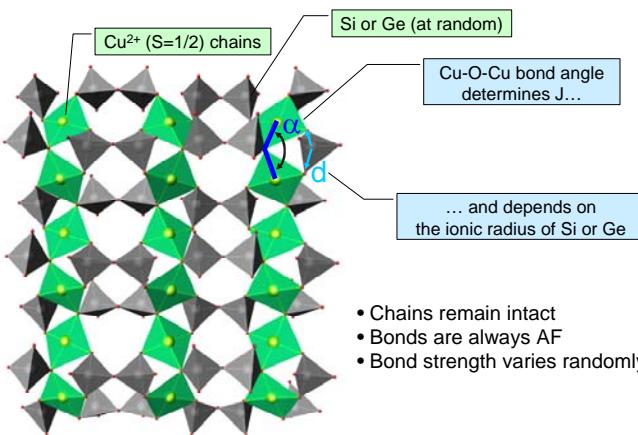


- Abundant weakly coupled large-scale singlets
- Divergent density of states at  $\omega \rightarrow 0$ :  $n(\omega) \sim 1/\omega \ln^3 \omega$
- Universal scaling laws:

$$\chi(T) = \frac{1}{T \ln^2(\Omega_0/T)} \quad S(q, \omega) = \frac{1}{\omega \ln^3(\Omega_0/\omega)} F(|q - q_0| d \frac{\delta J}{J} \ln^2(\Omega_0/\omega))$$

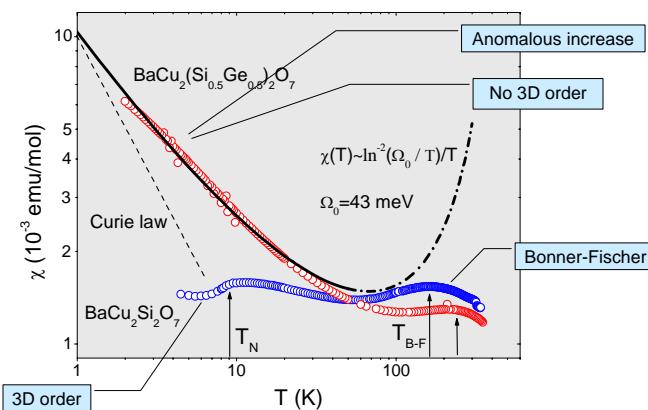
Doty, Fischer et al., 1992-1995  
K. Damle, O. Motrunich and D. Huse, PRL 84, 3434 (2000)

## A family of model compounds



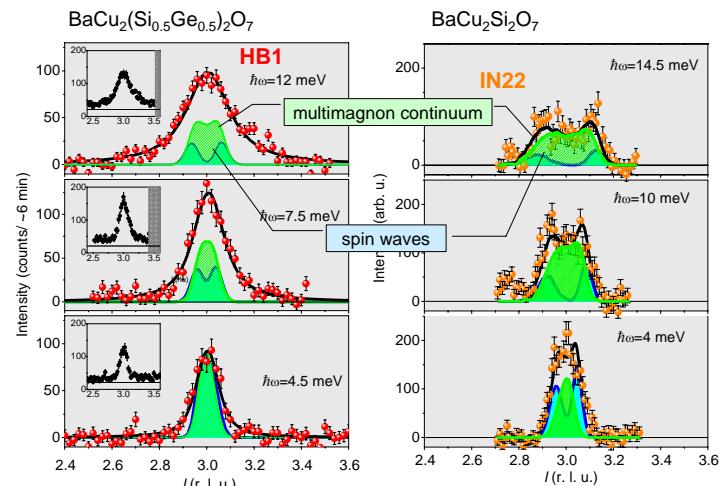
- $x=1$   $BaCu_2Si_2O_7$ ;  $J=24$  meV  $T_N=9.2$  K
- $x=0$   $BaCu_2Ge_2O_7$ ;  $J=50$  meV  $T_N=8.5$  K

## X=0.5: bulk properties



- $\chi(T)$  increase observed in all  $0 < x < 1$  compounds.
- Can be explained by divergent  $n(\omega)$  in random chains.

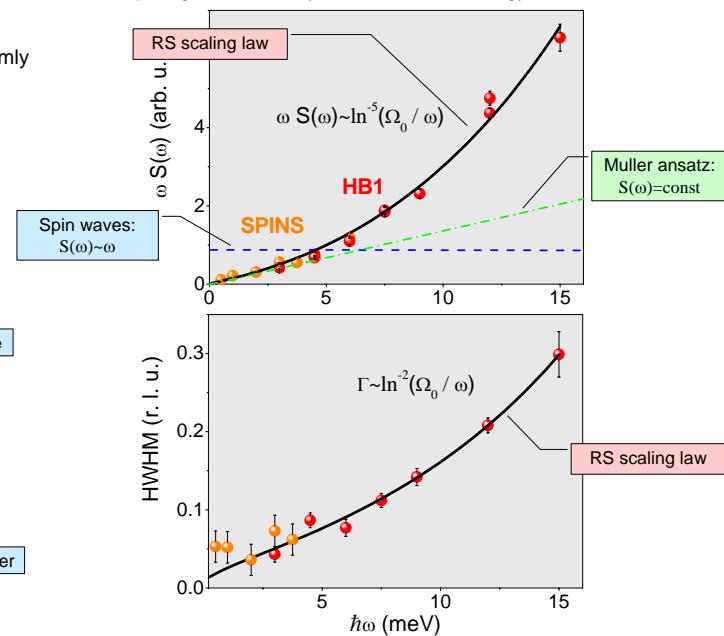
## Nature of spin excitations



- Broad Lorentzian shape
- Spread over entire zone
- **Diffusive** propagation of **local excitations**
- Width determined by size of excitations
- Typical "top hat" shape
- Confined to the  $\hbar\omega > |q|v$  cone
- **Coherent** propagation of **quasiparticles**
- Width determined by dispersion relation

## Scaling of the dynamic structure factor

- q-integrated intensity **increases** with energy!



- Scaling of HWHM allows to estimate  $\Delta J \sim 1$  meV

## Conclusion

**quenched disorder + quantum critical spin state**

=  
**universal scaling of spin dynamics**

